

Photovoltaic Devices

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Controllable Growth of Perovskite Films by Room-Temperature Air Exposure for Efficient Planar Heterojunction Photovoltaic Cells

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Abstract: A two-step solution processing approach has been established to grow void-free perovskite films for low-cost high-performance planar heterojunction photovoltaic devices. A high-temperature thermal annealing treatment was applied to drive the diffusion of CH₃NH₃I precursor molecules into a compact PbI2 layer to form perovskite films. However, thermal annealing for extended periods led to degraded device performance owing to the defects generated by decomposition of perovskite into PbI2. A controllable layer-by-layer spincoating method was used to grow "bilayer" CH3NH3I/PbI2 films, and then drive the interdiffusion between PbI2 and CH_3NH_3I layers by a simple air exposure at room temperature for making well-oriented, highly crystalline perovskite films without thermal annealing. This high degree of crystallinity resulted in a carrier diffusion length of ca. 800 nm and a high device efficiency of 15.6%, which is comparable to values reported for thermally annealed perovskite films.

Low-cost organometal halide perovskite solar cells have emerged as a prime candidate to meet the future energy generation demands on a gigawatt scale, owing to the certified power conversion efficiency (PCE) of around 20% with simple solution-based processing approaches.^[1-5] One-step solution-based approaches to grow hybrid perovskite thin films have typically generated a large density of pinholes or voids within perovskite films,^[6,7] which adversely impact the photovoltaic performance, especially in planar heterojunction configuration devices. Furthermore, large variations in film morphology significantly limited device performance reliabil-

ity, impacting yield. [8,9] As a result, two-step solution-based processing techniques such as methylammonium iodide (CH₃NH₃I) vapor assistance^[10] and the lead iodide (PbI₂)/ CH₃NH₃I bilayer interdiffusion.^[11–13] have been developed to grow void-free perovskite thin films.[10-14] These approaches employ high temperature (for example, ≥100°C) thermal annealing to drive the diffusion of CH₃NH₃I precursor molecules into the dense PbI2 layer to form compact methylammonium lead triiodide (CH₃NH₃PbI₃) perovskite thin films. However, thermal annealing for extended periods is known to cause the decomposition of perovskites into PbI₂, which degrades the device performance by acting as defects in the perovskite thin films.^[15] Therefore, finding simple and reliable solution processing approaches for void-free and highly crystalline perovskite films for efficient solar cells without thermal annealing is highly desirable for practical solar-cell manufacturing technology.

Herein we report a simple room-temperature, air-exposure process that removes the thermal annealing step and drives the interdiffusion of PbI₂ and CH₃NH₃I layers to synthesize highly crystalline and pinhole-free CH₃NH₃PbI₃ perovskite films. The interdiffusion process and the resulting formation of CH₃NH₃PbI₃ films were characterized by in situ X-ray diffraction (XRD), which revealed that ambient water molecules are a driving force for the interdiffusion. Perovskite thin-film photovoltaic devices resulting from the new air-exposure process displayed a PCE of 15.6%, which is comparable to that of most thermally annealed perovskite devices.

The PbI₂ layer and CH₃NH₃I layer were spin-coated sequentially onto TiO₂/ITO glass substrates using dimethylformamide (DMF) and 2-propanol, respectively. As shown in Figure 1a, the initial as-cast "bilayer" thin film not only exhibits Bragg peaks associated with PbI₂ crystal (001) planes at 12.54°, but also strong CH₃NH₃PbI₃ peaks at 14.05° and at 28.37° for the (110) and (220) planes, respectively, indicating the formation of perovskite crystals resulting from the CH₃NH₃I spin-coating process. It should be noted that another peak appeared at 13.89°, which was indexed as the perovskite CH₃NH₃PbI₃ (002) plane using Rietveld analysis (Supporting Information, Figure S1). Therefore, this initial stage of the as-cast film can be viewed as a "trilayer" film with a layer of CH₃NH₃PbI₃ perovskite crystals sandwiched between PbI₂ and CH₃NH₃I precursor layers (Figure 1b). The as-cast PbI₂/CH₃NH₃PbI₃/CH₃NH₃I "trilayer" films were exposed to ambient air at room temperature with a humidity of about 30%. After 60 min, the PbI₂ precursors were found to have completely converted into CH₃NH₃PbI₃ perovskites,

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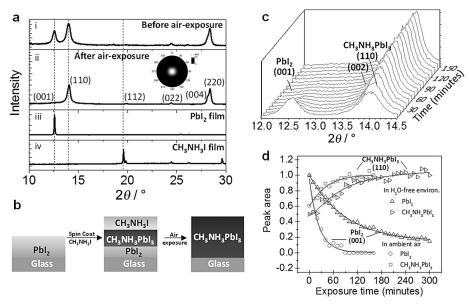


Figure 1. a) X-ray diffraction pattern of an as-cast glass/Pbl₂/CH₃NH₃I film before air exposure (i), after air-exposure for 60 min (ii), glass/Pbl₂ film (iii), and glass/CH₃NH₃I film (iv). Inset: X-ray pole figure acquired from the [110] direction for the air-exposure perovskite film. b) Depiction of the perovskite thin film growth process. c) Evolution of the X-ray diffraction pattern as a function of ambient air-exposure time by monitoring peaks at 12.54° (the (001) plane of Pbl₂), 13.89° (the (002) plane of CH₃NH₃Pbl₃), and 14.05° (the (110) plane of CH₃NH₃Pbl₃). d) A comparison of the Pbl₂ (001) plane and the CH₃NH₃Pbl₃ (110) plane peak areas as a function of exposure time for the air and water-free environment.

as indicated by the disappearance of the PbI₂ (001) Bragg peak (Figure 1a). X-ray pole figures were used to probe the global orientation and crystallographic texture of such films, and Figure 1a indicates that they were well-oriented with the CH₃NH₃PbI₃ (110) planes parallel to the surface of the substrate. Note that this plane has been theoretically identified as the most stable surface in the tetragonal phase of CH₃NH₃PbI₃.^[16] The SEM image shows that the CH₃NH₃PbI₃ perovskite film is free of pinholes, although a few large (ca. 1 μm) "flower-like" perovskite crystals are observed (Supporting Information, Figure S2). Atomic force microscopy was used to understand the surface roughness of the film and the large "flower-like" crystals. The root-mean-square roughness of a 20 μ m \times 20 μ m area on an air-exposed film was about 34 nm, and the large "flower-like" crystals (Supporting Information, Figure S2) have an average height of 70 nm, indicating they protrude above the surface of the film.

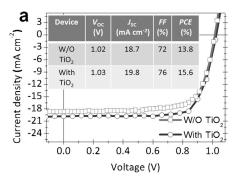
To investigate the evolution of the perovskite phase in the as-cast thin films during ambient exposure, in situ X-ray diffraction was employed to monitor the PbI₂ (001) and CH₃NH₃PbI₃ (110) Bragg peaks as a function of time. As shown in Figure 1 c, with increasing exposure time, the PbI₂ Bragg peak gradually decreased while the CH₃NH₃PbI₃ Bragg peak simultaneously increased (Figure 1 c,d), indicating the PbI₂ and CH₃NH₃I precursors were reacting to grow the CH₃NH₃PbI₃ crystals. After 60 min of exposure, the PbI₂ peak completely disappeared, and the intensity of the CH₃NH₃PbI₃ (110) peak correspondingly stabilizes, indicating that the reaction and crystallization were complete. To understand whether humidity was the driving force for the interdiffusion, air exposure studies were conducted on a similar initial PbI₂/

CH₃NH₃PbI₃/CH₃NH₃I "trilayer" film at room temperature in a water-free chamber (H₂O < 0.1 ppm). Interestingly, despite an essentially dry environment, the peak area intensity of the PbI₂ (001) plane was again found to decrease and the peak intensity of the CH₃NH₃PbI₃ (110) plane correspondingly increased, however more slowly (Supporting Information, Figure S3). This suggests that PbI₂ and CH₃NH₃I can gradually convert into CH3NH3PbI3 even in a water-free environment, although at a much lower rate as shown in Figure 1 d. It is thus hypothesized that ambient water accelerates the chemical reaction between PbI2 and CH₃NH₃I owing to the existence of a more reactive and metastable CH₃NH₃PbI₃·H₂O phase, which then spontaneously releases its water molecules at room temperature.[17,18]

To assess the applicability of the films for solar cells, prototype devices were fabricated. We first exam-

ined the photovoltaic performance of electron transporting layer (ETL)-free devices (ITO/CH₃NH₃PbI₃/2,2',7,7'-tetrakis-(N,N-di-pmethoxyphenylamine)-9,9'-spirobifluorene (Spiro-OMeTAD)/Ag) based on our air-exposed perovskite films. To grow compact perovskite layers, a UV-ozone treatment for 10 min was used to generate a hydrophilic ITO glass surface, as shown by water contact-angle measurements (Supporting Information, Figure S4). Surprisingly, when swept from forward bias to reverse bias, the device without ${\rm TiO_2}$ yielded a $J_{\rm SC}$ of 18.7 mA cm⁻², a $V_{\rm OC}$ of 1.02 V, a FF of 72%, and a PCE of 13.8% (Figure 2a). The cross-sectional SEM image acquired from such a high-performance TiO₂-free device shows a compact perovskite layer (Supporting Information, Figure S5), which is consistent with previous results.^[7] To verify the $J_{\rm SC}$ from the J-V curve, the external quantum efficiencies (EQEs) were measured (Supporting Information, Figure S6). By integrating the EQE curve across the standard AM 1.5 G solar spectrum (100 mA cm⁻²), a J_{SC} of 18.2 mA cm⁻² was calculated, which is consistent with the J_{SC} measured from the J-V scan. For comparison, a TiO₂ ETL was incorporated into the devices, and as shown in Figure 2a, the typical device exhibited improved performance; that is, a $J_{\rm SC}$ of 19.8 mA cm⁻², a $V_{\rm OC}$ of 1.03 V, a high FF of 76 %, and a PCE of 15.6% when swept from forward bias to reverse bias. It should be noted that both types of devices, with and without TiO2, exhibited hysteresis characteristics as shown in the Supporting Information, Figure S7. To reveal how the J-Vhysteresis affects the maximum power output from our devices, we measured a device in both forward and reverse scan mode, which showed PCEs of 15.2% and 8%, respectively (Supporting Information, Figure S7). We then acquired





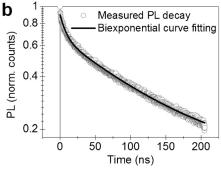


Figure 2. a) Comparison of *J*–*V* curves of devices without (□) and with (○) TiO_2 acquired under standard illumination (AM 1.5 G, 100 mW cm⁻²). The obtained photovoltaic parameters are shown in the inset table. b) Time-resolved photoluminescence (PL) at about 760 nm (○), showing a biexponential decay characteristic (——).

the photocurrent as a function of time measured at the maximum power output point (0.792 V) under an illumination of $100 \,\mathrm{mW \, cm^{-2}}$. The photocurrent density saturated at about $16 \,\mathrm{mA \, cm^{-2}}$, and the device showed a stable PCE of about $12.7 \,\%$, despite the strong J-V hysteresis.

To further understand the origin of their high photovoltaic performance, we first characterized the time-resolved photoluminescence (PL) at about 760 nm (Figure 2b). Based on the biexponential fitting results, two lifetimes were obtained: a short lifetime of 8.3 ns and a long lifetime of 114.2 ns. The fast decay is ascribed to significant surface recombination, while the slow decay is attributed to the dominant free carrier decay within the bulk film.^[19] Hall measurements of the airexposed perovskite films showed a carrier mobility of about 2.16 cm²V⁻¹s⁻¹ (Supporting Information, Figure S8), and when combined with carrier lifetime measured previously, the carrier diffusion length was estimated to be about 800 nm, which is comparable to the carrier diffusion length in most thermally annealed CH₃NH₃PbI₃ perovskite thin films.^[13] Therefore, the high photovoltaic performance for both types of air-exposed perovskite film devices (with and without TiO₂) should be partially attributed to this long carrier diffusion length.

To understand the cross-sectional morphology and crystallinity of the devices, cross-sectional TEM images and SAED studies were performed. In Figure 3a, the compact CH₃NH₃PbI₃ perovskite layer was clearly identified in the cross-sectional TEM image of our high performance device, which has an ITO/TiO₂/CH₃NH₃PbI₃/Spiro-OMeTAD/Ag architecture (Figure 3a). A series of SAED patterns were

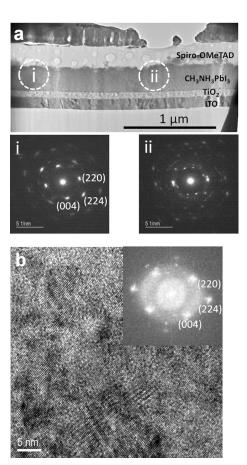


Figure 3. a) Bright-field cross-sectional TEM image of an entire device with a $\rm TiO_2$ layer. Two typical SAED patterns (i and ii) were acquired from $\rm CH_3NH_3PbI_3$ layer. Pattern i shows identical crystallographic orientation, whereas pattern ii shows multiple grain orientations within the aperture diameter size of 325 nm. b) High-resolution TEM image taken from the $\rm CH_3NH_3PbI_3$ layer. Inset: the corresponding FFT pattern, viewed along the [110] zone axis. Note that indexing of the (004) and (220) diffraction spots could possibly be switched, because the similar d spacing prevents unambiguous determination of these spot locations.

acquired along the direction parallel to the glass substrate to examine the crystallographic features of the perovskite film. Two typical SAED patterns are shown in Figure 3a. We observed a polycrystalline spot pattern (pattern i) with each domain oriented very similarly over a distance of 325 nm (the aperture diameter), suggesting a preferential growth direction, which is consistent with a preferred crystal orientation of [110] as revealed by X-ray pole figure measurements. Pattern i is consistent with previously reported observations where CH₃NH₃PbI₃ is viewed along the [110] zone axis.^[20] However, it should be noted that we also observed regions where the crystals did not orient similarly, as shown by pattern ii. A high-resolution TEM image (Figure 3b) and the corresponding fast Fourier transform (FFT; inset of Figure 3b) were taken from the ambient-air-exposed perovskite layer, and clear lattice fringes and a well-ordered crystal lattice were observed, despite the presence of surface damage owing to focused-ion beam milling during sample preparation. Thus, it is concluded that the high crystallinity and highly



oriented nature of the perovskite layer contributes significantly to the long carrier diffusion in ambient-air-exposure perovskite thin films.

Electron-beam-induced current (EBIC) measurements were employed to further investigate the device operation mechanism. The cross-sectional SEM and corresponding EBIC images are shown in Figure 4a. The line profiles of the SEM and corresponding EBIC were taken across the

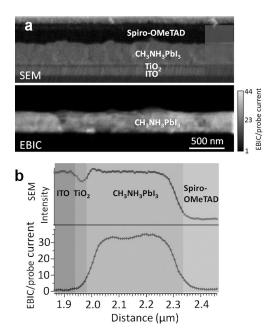


Figure 4. a) SEM (top) and EBIC (bottom) images of a device with a TiO₂ layer, with the corresponding average line profiles shown in (b). The line profiles in (b) were taken in the region displayed in (a).

entire device, as indicated in Figure 4b. The EBIC signal was only observed within the CH₃NH₃PbI₃ perovskite layer, indicating that the perovskite layer is the only layer generating free charge carriers that can be collected by the electrodes. A uniform carrier generation was observed in the perovskite film with slight peaks close to the perovskite/TiO2 and perovskite/Spiro-OMeTAD inherent for a p-i-n structure.

In summary, a simple room-temperature air-exposure approach was developed to grow highly crystalline, welloriented perovskite films, resulting in a long carrier diffusion length of about 800 nm. The devices with and without TiO₂ ETL yielded PCEs of 15.6% and 13.8%, respectively, which are comparable to those of most thermally annealed perovskite devices. The non-thermal-annealing, room-temperature air-exposure approach is compatible with roll-to-roll processing, which enables very low-cost flexible photovoltaics.

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Keywords: air exposure · in situ X-ray diffraction · perovskites · photovoltaic devices · thin films

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- [1] J. Burschka, N. Pellet, S.-J. Moon, R. Humphry-Baker, P. Gao, M. K. Nazeeruddin, M. Gratzel, Nature 2013, 499, 316-319.
- [2] S. Das, B. Yang, G. Gu, P. C. Joshi, I. N. Ivanov, C. M. Rouleau, T. Aytug, D. B. Geohegan, K. Xiao, ACS Photonics 2015, 2, 680-
- [3] F. Hao, C. C. Stoumpos, D. H. Cao, R. P. H. Chang, M. G. Kanatzidis, Nat. Photonics 2014, 8, 489-494.
- [4] N. J. Jeon, J. H. Noh, W. S. Yang, Y. C. Kim, S. Ryu, J. Seo, S. I. Seok, Nature 2015, 517, 476-480.
- [5] B. Yang, O. Dyck, J. Poplawsky, J. Keum, A. Puretzky, S. Das, I. N. Ivanov, C. Rouleau, G. Duscher, D. B. Geohegan, K. Xiao, J. Am. Chem. Soc. 2015, 137, 9210-9213.
- [6] G. E. Eperon, V. M. Burlakov, P. Docampo, A. Goriely, H. J. Snaith, Adv. Funct. Mater. 2014, 24, 151-157.
- W. Ke, G. Fang, J. Wan, H. Tao, Q. Liu, L. Xiong, P. Qin, J. Wang, H. Lei, G. Yang, M. Qin, X. Zhao, Y. Yan, Nat. Commun. 2015, 6, DOI: 10.1038/ncomms7700.
- [8] B. Conings, L. Baeten, C. De Dobbelaere, J. D'Haen, J. Manca, H.-G. Boyen, Adv. Mater. 2014, 26, 2041 - 2046.
- [9] W. Zhang, M. Saliba, D. T. Moore, S. K. Pathak, M. T. Hörantner, T. Stergiopoulos, S. D. Stranks, G. E. Eperon, J. A. Alexander-Webber, A. Abate, A. Sadhanala, S. Yao, Y. Chen, R. H. Friend, L. A. Estroff, U. Wiesner, H. J. Snaith, Nat. Commun. 2015, 6, DOI: 10.1038/ncomms7142.
- [10] Q. Chen, H. Zhou, Z. Hong, S. Luo, H.-S. Duan, H.-H. Wang, Y. Liu, G. Li, Y. Yang, J. Am. Chem. Soc. 2014, 136, 622-625.
- [11] C.-W. Chen, H.-W. Kang, S.-Y. Hsiao, P.-F. Yang, K.-M. Chiang, H.-W. Lin, Adv. Mater. 2014, 26, 6647-6652.
- [12] Y. Chen, T. Chen, L. Dai, Adv. Mater. 2015, 27, 1053-1059.
- [13] Z. Xiao, Q. Dong, C. Bi, Y. Shao, Y. Yuan, J. Huang, Adv. Mater. **2014**, 26, 6503 - 6509.
- [14] J.-H. Im, I.-H. Jang, N. Pellet, M. Grätzel, N.-G. Park, Nat. Nanotechnol. 2014, 9, 927-932.
- [15] C. Bi, Y. Shao, Y. Yuan, Z. Xiao, C. Wang, Y. Gao, J. Huang, J. Mater. Chem. A 2014, 2, 18508-18514.
- [16] J. Haruyama, K. Sodeyama, L. Han, Y. Tateyama, J. Phys. Chem. Lett. 2014, 5, 2903-2909.
- [17] F. Hao, C. C. Stoumpos, Z. Liu, R. P. Chang, M. G. Kanatzidis, J.
- Am. Chem. Soc. 2014, 136, 16411 16419. [18] H. Zhou, Q. Chen, G. Li, S. Luo, T.-b. Song, H.-S. Duan, Z.
- Hong, J. You, Y. Liu, Y. Yang, Science 2014, 345, 542-546. D. Shi, V. Adinolfi, R. Comin, M. Yuan, E. Alarousu, A. Buin, Y. Chen, S. Hoogland, A. Rothenberger, K. Katsiev, Y. Losovyj, X. Zhang, P. A. Dowben, O. F. Mohammed, E. H. Sargent, O. M.
- Bakr, Science 2015, 347, 519-522. [20] H. Zhu, Y. Fu, F. Meng, X. Wu, Z. Gong, Q. Ding, M. V. Gustafsson, M. T. Trinh, S. Jin, X. Zhu, Nat. Mater. 2015, 14, 636 - 642.

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